Dynamical aspects and the role of IVR for the reactivity of noble metal clusters towards molecular oxygen

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Abstract. Here, we present the dynamical aspects and the role of internal energy redistribution (IVR) in the reactivity of noble metal clusters towards O_2 . We show on the example of Ag₃O₂ / Ag₃O₂ / Ag_3O_2^+ that NeNePo spectroscopy carried out under zero electron kinetic energy (ZEKE) conditions can be a powerful tool to investigate the geometry relaxation and IVR induced by photodetachment in real time. Furthermore, we demonstrate that difference in the reactivity of Ag_{^{*6*}</sub></sub> and Au₆[−] towards O₂ can be</sub>} attributed to different nature of the IVR process. Dissipative IVR in Ag*[−]* ⁶ favors fast complex stabilization, whereas resonant IVR found for Au*[−]* ⁶ might be an important factor determining the catalytic activity of Au*[−]* ⁶ cluster in the CO oxidation.

PACS. 31.15.Qg Molecular dynamics and other numerical methods – 31.15.Ar Ab initio calculations

1 Introduction

Noble metal clusters exhibit fascinating reactive properties such as strongly size and charge dependent reactivity towards small molecules such as O_2 , CO etc. [1–3]. Moreover, gold clusters have attracted significant attention due to their ability to catalyze the oxidation of CO at low temperatures [4–6]. In this context, the understanding of the mechanism for the activation of molecular oxygen is of fundamental importance. This has stimulated a number of theoretical studies in which specific geometric as well as electronic properties which are responsible for cooperative adsorption of multiple adsorbate molecules on small noble metal clusters, have been identified recently [2,6,7]. In particular, the cooperative effects are regarded essential to the catalytic activity of the gas-phase noble metal clusters in the CO combustion reaction [8,9]. Moreover, since both experimental and theoretical studies have shown that the clusters alone are not capable to break the O–O bond [1,10], the question can be raised whether exploration of reactivity exclusively based on the stationary energetics is sufficient or whether dynamical processes, e.g. internal energy redistribution (IVR) [11], might become the key issue for full understanding of the catalytic properties.

Here, we wish to address two complementary aspects of reaction dynamics of silver and gold clusters towards O_2 . One is to study reaction dynamics in neutral clusters in the frame of the negative-to-neutral-to-positive (NeNePo) spectroscopy [12–16]. After an initial ensemble is generated in the anionic electronic ground state, photodetachment of the excess electron occurs and subsequently, dynamics in the neutral state is monitored by probing the cationic state with a laser pulse of sufficiently high energy for ionization. The goal of theoretical fs-NeNePo spectroscopy is to provide conditions under which different processes, such as bond breaking/formation, isomerization and IVR, and their time scales can be observed [17]. The potential prospects of the NeNePo spectroscopy for real time laser spectroscopic investigations of the activation of molecular $O₂$ will be illustrated on the prototype example of Ag_3O_2 .

The second aspect we would like to address concerns differences in reactivity of gold and silver clusters towards O² which might be caused by inherently different dynamical properties. It has been found experimentally that Au6O*[−]* ² is catalyzing the oxidation of CO [18] in contrast to Ag6O*[−]* ² [19] which is inactive. However, structural and binding properties of Ag₆O₂[−] and Au₆O₂[−] do not differ substantially. Therefore, the question can be raised whether the pronounced relativistic effects in gold [20–22] might be responsible for the found catalytic activity by invoking a special type of IVR which is reflected in the reaction dynamics and energy redistribution. Therefore, we investigated the differences in the reaction dynamics between silver and gold clusters by studying the collision dynamics of Ag₆[−] and Au₆[−] with O₂. In particular, we will highlight the role of resonant vs. dissipative IVR in the stabilization of complexes with O_2 .

2 Computational

The time-resolved NeNePo spectra were simulated in the framework of our Wigner distribution approach [13] based

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Fig. 1. Scheme of the multistate fs-dynamics for $Ag_3O_2^-/Ag_3O_2/Ag_3O_2^+$ in the framework of NeNePo pumpprobe-spectroscopy with structures and energy intervals (in eV) for pump and probe steps.

on the propagation of an ensemble of classical trajectories carried out "on the fly". Because here the dynamics in the ground electronic states of the neutral species is required, we employ the gradient corrected density functional theory with the atomic orbitals (Gaussian basis sets). For details on electronic structure calculations, cf. reference [17]. The collision dynamics of Ag₆[–] and Au₆[–] with O₂ has been investigated employing DFT methods with the gradient corrected PBE functional [23] using Stuttgart 19e*[−]* RECP [24] for Ag and Au and triple zeta polarization basis set for O [25]. The collision parameters have been sampled resulting in ≈ 400 initial conditions. Collisional impact energy of 0.25 eV has been chosen corresponding to room temperature.

3 Results and discussion

*Simulation of NeNePo spectra for Ag*3*O[−]* ² */ Ag*3*O*² */ Ag*3*O*⁺ 2

The electronic and structural aspects of adsorption of molecular oxygen on anionic silver clusters have been investigated previously [7]. Here, we present the simulation of the NeNePo spectra of Ag₃O₂ / Ag₃O₂ / Ag₃O₂ [16] with the specific goal to address the energy redistribution process which occurs upon geometry relaxation induced by the photodetachment process.

In Figure 1, the scheme for the multistate dynamics of the Ag_3O_2 ⁻ / Ag_3O_2 / Ag_3O_2 ⁺ system is shown. In the anionic state, two isomers with linear (A I) and triangular (A II) Ag₃ subunit are very close in energy (0.179 eV) . The exact energy sequence is very difficult to determine, even with highly accurate methods. On the other hand, the difference in the vertical detachment energies (VDE) between the two isomers (2.8 eV for isomer A I and 2.2 eV for isomer A II) allows for selective excitation of only one

Fig. 2. Bunch of time dependent cation-neutral energy gaps of Ag3O2. Trajectories leading to mixture of isomers II and III are marked by light color and trajectories leading to isomer I are marked by black color. Horizontal black lines indicate energies for the probe-pulses at 6.8, 7.3 and 7.6 eV.

isomer by choice of the appropriate pump-pulse energy. Therefore, in the simulations presented here we assume that only the linear isomer A I is populated. A canonical Wigner distribution function is used to generate the initial conditions at $T = 20$ K.

After photodetachment with a pump-pulse energy of 2.75 eV, a transition state on the neutral PES of Ag_3O_2 is reached. As shown in Figure 1, three isomers are present in the neutral state which mainly differ in the binding situation of the O_2 subunit. In the global minimum (I) oxygen bridges two silver atoms, whereas in the two higher lying isomers oxygen is bound to one silver atom, forming two (II) or one bond (III), respectively (cf. Fig. 1).

The dynamics in the neutral state induced by the photodetachment is dominated by relaxation of the Ag₃ subunit from the linear to the triangular form. This process is completed within 1 ps, as can be seen from the time dependent energy gaps between the neutral and the cationic state presented in Figure 2 (light lines). Subsequently, the energy gained during this process is partially transferred to the molecular oxygen, leading to a mixture of neutral isomers II and III. At later times, two different isomerization pathways can be identified from Figure 2. A mixture of the two neutral isomers II and III with higher energies still remains dominant channel (light lines, 90%) while the channel leading to the global minimum (I) amounts only to about 10% of the population within the simulation period of 5 ps (cf. black curves in Fig. 2).

The simulated signal under zero electron kinetic energy condition (NeNePo-ZEKE) and including continuum (NeNePo) for three different probe-pulse energies 6.8, 7.3 and 7.6 eV (cf. horizontal lines in Fig. 2) are presented in Figure 3. For the probe-pulse energy of 6.8 eV, the NeNePo-ZEKE signal reflects the main isomerization process into a mixture of the two higher lying neutral isomers due to close lying values of the ionization potentials of 6.8 eV (isomer III) and 7.18 eV (isomer II). Since the contribution of the continuum for the 6.8 eV probe-pulse energy is very small, there is no significant difference between the NeNePo and NeNePo-ZEKE signals.

Under ZEKE conditions, the probe-pulse energy of 7.3 eV allows monitoring the onset of the relaxation dynamics of the Ag₃ subunit at ≈ 500 fs. However, if the

Fig. 3. Simulated NeNePo-ZEKE and NeNePo signals of Ag3O² for 20 K initial condition ensemble at probe-energies 6.8, 7.3 and 7.6 eV.

continuum is included the signal increases for 800 fs and then remains constant which means that relaxation dynamics is smeared out. The ZEKE signal with a probepulse energy of 7.6 eV exhibits a fast decay within 1 ps due to the escape from the Franck-Condon region, which is initially populated after the photodetachment. Moreover, after 1.5 ps the signal starts to rise again indicating the increasing population of the global minimum (IP = 7.62 eV). This signal offers the opportunity to identify isomerization processes leading to the global minimum. The global minimum starts to be populated at 1.5 ps and the population reaches $\approx 10\%$ after 4 ps. On the contrary, the NeNePo signal for the same energy does not allow to distinguish any processes due to the fact that all structures are ionized.

In order to study the IVR process initiated by the photodetachment we present in Figure 4 the distribution of the kinetic energy between the $Ag₃$ and the $O₂$ fragments. The oscillations of the kinetic energy of O_2 within the first 700 fs reflects the thermally excited vibrations of the molecular oxygen. For the Ag³ subunit, the kinetic energy starts to rise due to relaxation processes up to a maximum at 850 fs, which corresponds to the time of the closest approach of both silver atoms (internal collision). The timescale for this relaxation process is quite similar to those of the free $Ag₃$ [13], in which internal collision occurs after 950 fs. Subsequently, the increasing distance of the silver atoms after the collision results in a decay of the kinetic energy up to a minimum at 1270 fs. The parallel increase of the O_2 kinetic energy with a maximum at 950 fs indicates partial energy transfer from the Ag³ to the O_2 fragment. A second maximum of the Ag_3 kinetic energy at 1450 fs can be assigned to the second collision of the silver atoms, also followed by a partial energy transfer to the oxygen fragment.

The resulting increase of kinetic energy in the O_2 fragment leads to the two different isomerization pathways described above. The kinetic energy in the oxygen fragment is not sufficient to overcome the energy barrier for the isomerization into the global minimum. Therefore, only about 10% of the ensemble is able to reach the global

Fig. 4. Distribution of the kinetic energy between Ag³ and O² fragment.

minimum within the simulated time interval. The most efficient isomerization occurs between isomer II and III.

$Collisions of Ag₆[−] and Au₆[−] with O₂$

In order to address the influence of the dynamics and in particular the IVR process on the reaction between Ag*−* 6 and Au_6^- with molecular O_2 , we compared the reaction dynamics by simulating collisions between clusters and O2. In particular, the question, whether different reaction dynamics caused by the relativistic effects which are more pronounced for gold than for silver might be responsible for the observed different catalytic activity will be addressed.

As can be seen from Figure 5, in the case of Ag_{6}^- + O_{2} , dissipative IVR from O_2 to cluster can be identified. The kinetic energy of the complex is mainly localized in the Ag⁶ subunit (black curve in Fig. 5a) after the first 1.4 ps. This efficient process leads to fast stabilization of the complex and furthermore, prevents activation of O_2 . Moreover, due to excess of adsorption energy, a large deviation from planarity is observed as can be seen from the snapshots of the MD trajectory (inset in Fig. 5a). In contrast, in the case of Au_6^- + O_2 a resonant IVR is present which can be seen from Figure 5b. The kinetic energy is periodically distributed over both Au_6 and O_2 subgroups for a longer period of time allowing $O₂$ to remain activated for longer times. In contrast to Ag₆⁻, no large deviation from planarity has been identified. Consequently, in the case of gold clusters the adsorption energy remains localized at the adsorption site and therefore might be used to more easily overcome reaction barriers.

Notice that the differences in energy redistribution are not caused by different masses but are present due to inherent different dynamical processes connected with higher rigidity of Au*[−]* ⁶ caused by pronounced relativistic effects. This leads to more directional bonding in Au*[−]* ⁶ due to the increased contribution of d -electrons [21,26]. This directionality is also reflected in more rigid structure of Au₆[−]. In contrast, the s-like delocalized bonding in $Ag₆$ [−] gives rise to flat energy surface with close lying isomers.

In order to prove the higher rigidity of Au*[−]* ⁶ , MD simulations of the isomerization dynamics of Ag₆[–] and Au*[−]* ⁶ have been performed at different temperatures (cf. Fig. 5c). We use Lindemann index δ as a measure of rigidity. At similar temperatures δ has much higher values for silver than for gold indicating higher rigidity of Au₆[−]. Clearly, the Lindemann index δ reaches higher values at

Fig. 5. (a, b) Energy partition between Ag₆ (Au₆) and O₂ for two reactive collisions demonstrating dissipative IVR for Ag[−] (a) and resonant IVR for Au[−] (b) For illustration representative single trajectory h (a) and resonant IVR for Au_o[−] (b). For illustration, representative single trajectory has been chosen. Red: $E_{kin}(O_2)$; black: $E_{kin}(Ag_6/Au_6)$. (c) Lindemann indices δ for Ag_6^- (dashed line) and Au_6^- (solid line) at different temperatures.

lower temperature for silver compared to gold indicating melting at lower temperature in the former case.

In summary, we have shown that the energy redistribution process is much more efficient in the case of Ag*−* 6 than for Au*−* ⁶ . This leads to the fast stabilization of the complex with O_2 . In contrast to that, upon collision of Au*−* ⁶ with O² resonant IVR process occurs which might facilitate the overcoming of the barriers in subsequent reactions with CO. Thus we have shown that dynamical aspects might significantly influence the reactivity properties of clusters and can be an important factor which determines potential catalytic activity.

4 Conclusion

We have presented two complementary theoretical investigations which allowed us to study the energy redistribution in neutral and anionic complexes between noble metal clusters and molecular oxygen. We have shown that geometric relaxation and IVR process can be identified in real time in the framework of the NeNePo spectroscopy. Our theoretical simulation predicts, that it is mandatory to carry out experiments close to ZEKE condition in order to distinguish different processes in the NeNePo spectra. Furthermore, we have demonstrated that different IVR processes can significantly influence the reactivity of anionic gold and silver clusters towards O_2 . By studying collisional dynamics of Ag_{6}^- and Au_{6}^- with O_{2} we have shown that differences in their reactivity can be attributed to different nature of the IVR process. Fast dissipative IVR in Ag[−] is favorable for the complex stabilization leading to a complex which is not catalytically active. In contrast, the resonant IVR process in Au*[−]* ⁶ gives rise to an energy localization at the active site which might facilitate O–O bond breaking in the reaction with CO.

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